

**9. Exercise Sheet, Lecture Theoretical Quantum and Atom Optics**

University of Hamburg, Prof. P. Schmelcher

To be returned on 09.06.2011, in the tutorials

**Exercise 16.** Landau-Zener formula

- (a) Consider a two-level system described by the Hamiltonian

$$H = \begin{pmatrix} -\epsilon & -f \\ -f & \epsilon \end{pmatrix}, \quad (1)$$

with a linearly time dependent sweep  $\epsilon = -\alpha t$ , where  $f, \alpha > 0$ .

Show that the Schrödinger equation (setting  $\hbar \equiv 1$ )

$$i \frac{\partial}{\partial t} \begin{pmatrix} c_1(t) \\ c_2(t) \end{pmatrix} = H \begin{pmatrix} c_1(t) \\ c_2(t) \end{pmatrix} \quad (2)$$

leads to the ordinary differential equation

$$\frac{d^2}{dt^2} c_2(t) + [f^2 - i\alpha + (\alpha t)^2] c_2(t) = 0 \quad (3)$$

for the amplitude  $c_2(t)$ .

- (b) Show that the variable substitution  $t \rightarrow z(t) = e^{-i\frac{\pi}{4}} (2\alpha)^{\frac{1}{2}} t$  transforms Eq. (3) to the so called Weber equation,

$$\frac{d^2}{dz^2} \tilde{c}_2(z) + \left[ \nu + \frac{1}{2} - \frac{1}{4} z^2 \right] \tilde{c}_2(z) = 0 \quad (4)$$

with  $\nu = \frac{if^2}{2\alpha}$ , where  $\tilde{c}_2(z) = c_2(z(t))$ .

- (c) Eq. (4) is solved by the four parabolic cylinder (Weber) functions  $D_\nu(z)$ ,  $D_\nu(-z)$ ,  $D_{-\nu-1}(iz)$  and  $D_{-\nu-1}(-iz)$ , the former two of which are linearly independent for  $\nu \notin \mathbb{Z}$ .

Verify that  $D_\mu(\zeta)$  is a solution of the Weber equation if it obeys the following recursive relations for arbitrary  $\zeta$  and  $\mu$ :

$$D_{\mu+1}(\zeta) - \zeta D_\mu(\zeta) + \mu D_{\mu-1}(\zeta) = 0, \quad (5)$$

$$\frac{d}{d\zeta} D_\mu(\zeta) + \frac{1}{2} \zeta D_\mu(\zeta) - \mu D_{\mu-1}(\zeta) = 0 \quad (6)$$

- (d) Assume that, initially, the system is in the state  $|1\rangle$ , so that the following initial conditions hold:

$$|c_1(t \rightarrow -\infty)|^2 = 1, \quad (7)$$

$$|c_2(t \rightarrow -\infty)|^2 = 0. \quad (8)$$

During the linear sweep over the avoided crossing, the coupling  $f$  causes population transfer from  $|1\rangle$  to  $|2\rangle$ , and the aim is to find the final distribution

$$P_{LZ} \equiv |c_1(t \rightarrow \infty)|^2 = 1 - |c_2(t \rightarrow \infty)|^2. \quad (9)$$

Among the four solutions of Eq. (4) only the Weber function  $D_{-\nu-1}(-iz(t))$  vanishes for  $t \rightarrow -\infty$ . The amplitude  $c_2(t)$  thereby fulfills the initial condition 8 if it is written as

$$c_2(t) = \tilde{c}_2(z(t)) = AD_{-\nu-1}(-iz(t)). \quad (10)$$

where  $A$  is a normalization constant. Defining  $R \equiv \sqrt{2\alpha t}$ , the asymptotic expressions for  $D_{-\nu-1}(-iz(t))$  in the two limits  $t \rightarrow \mp\infty$  are given by

$$D_{-\nu-1}(-iz(t \rightarrow -\infty)) = e^{-\frac{1}{4}\pi(\nu+1)i} e^{-i\frac{R^2}{4}} R^{-\nu-1} \quad (11)$$

$$D_{-\nu-1}(-iz(t \rightarrow +\infty)) = \frac{\sqrt{2\pi}}{\Gamma(\nu+1)} e^{\frac{1}{4}\pi\nu i} e^{i\frac{R^2}{4}} R^\nu \quad (12)$$

(i) Find the asymptotic expression  $c_1(t \rightarrow -\infty)$ , by substituting  $c_2(t \rightarrow -\infty)$  into the Schrödinger equation (2).

*Hint:* Use the chain rule  $\frac{dc_2}{dt} = \frac{dc_2}{dR} \frac{dR}{dt}$ .

(ii) Show that  $A = \sqrt{\gamma} e^{-\frac{\pi\gamma}{4}}$  fulfills the normalization condition (7), where  $\gamma = -i\nu = \frac{f^2}{2\alpha}$ .

(iii) Derive the Landau-Zener formula:

$$P_{LZ} = e^{-2\pi\gamma}. \quad (13)$$

*Hint:* Calculate  $|c_2(t \rightarrow \infty)|^2$  using the properties of the Gamma function:

$$\Gamma(\pm i\gamma + 1) = \pm i\gamma \Gamma(\pm i\gamma) \quad (14)$$

$$|\Gamma(\pm i\gamma)| = \sqrt{\frac{\pi}{\gamma \sinh \pi\gamma}} \quad (15)$$

**10 Points**