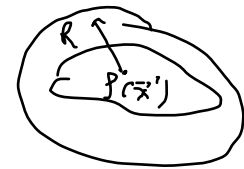


## 18. Multipolentwicklung

Gegeben: beliebig, lokalisiert Ladungsverteilung  $\rho(\vec{x}')$

$$\text{mit } \rho(\vec{x}') = \begin{cases} \text{beliebig} & \text{für } |\vec{x}'| < R \\ 0 & \text{für } |\vec{x}'| \geq R \end{cases}$$



und keine weitere Leiter im Raum  $\Rightarrow$  keine R.D.

$$\Rightarrow \Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int_{|\vec{x}'| < R} \frac{\rho(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x', \quad \Delta\Phi = -\frac{\rho}{\epsilon_0}$$

Wenn Integral analytisch mit behandelt, dann  
das Fernfeld approximativ bestimmt werden.

$$\Phi(\vec{x}) \text{ für } |\vec{x}| \gg R$$

Multipolentwicklung = Entwicklung in  $\frac{\vec{x}'}{|\vec{x}|} \ll 1$

$$\begin{aligned} \frac{1}{|\vec{x}-\vec{x}'|} &= \left( (x-x')^2 + (y-y')^2 + (z-z')^2 \right)^{-\frac{1}{2}} = \left( x^2 \left(1 - \frac{x'}{x}\right)^2 + y^2 \left(1 - \frac{y'}{y}\right)^2 + z^2 \left(1 - \frac{z'}{z}\right)^2 \right)^{-\frac{1}{2}} \\ &= \left( \sum_{i=1}^3 X_i^2 (1-a_i)^2 \right)^{-\frac{1}{2}}, \quad a_i = \frac{x'_i}{X_i}, \quad X_1 = x, X_2 = y, X_3 = z \end{aligned}$$

Taylor Entwicklung in  $a_i$

$$\begin{aligned} \text{allgemein: } f(a_i) &= f(0) + \sum_{i=1}^3 a_i \left. \frac{\partial f}{\partial a_i} \right|_{a_i=0} + \frac{1}{2} \sum_{i,j} a_i a_j \left. \left( \frac{\partial}{\partial a_i} \frac{\partial}{\partial a_j} f \right) \right|_{a_i=0} \\ &\quad + O(a^3) \end{aligned}$$

$$\left. \frac{1}{|\vec{x}-\vec{x}'|} \right|_{a_i=0} = \frac{1}{r}$$

$$\frac{\partial}{\partial a_i} \frac{1}{|\vec{X} - \vec{X}'|} \Big|_{a_i=0} = \frac{X_i^2 (1-a_i)}{\left(\sum_k X_k^2 (1-a_k)^2\right)^{3/2}} \Big|_{a_i=0} = \frac{X_i^2}{r^3}$$

$$\begin{aligned} \frac{\partial}{\partial a_j} \frac{\partial}{\partial a_i} \frac{1}{|\vec{X} - \vec{X}'|} \Big|_{a_i=0} &= -\frac{X_i^2 \delta_{ij}}{r^3} + \frac{3}{r^5} \frac{X_i^2 (1-a_i)}{\left(\sum_k X_k^2 (1-a_k)^2\right)^{3/2}} \cdot (2X_j^2 (1-a_j)) \Big|_{a_j=0} \\ &= -\frac{X_i^2 \delta_{ij}}{r^3} + 3 \frac{X_i^2 X_j^2}{r^5} \end{aligned}$$

$$\begin{aligned} \frac{1}{|\vec{X} - \vec{X}'|} &= \frac{1}{r} + \sum_i a_i \frac{X_i^2}{r^3} + \frac{1}{2} \sum_i \sum_j a_i a_j \left( \frac{3X_i^2 X_j^2}{r^5} - \frac{X_i^2 \delta_{ij}}{r^3} \right) + O(a^3) \\ &\stackrel{a_i = \frac{X_i}{r}}{\Downarrow} = \frac{1}{r} + \sum_i \frac{X_i^1 X_i^1}{r^3} + \frac{1}{2} \sum_i \sum_j \left( \frac{3X_i^1 X_j^1 X_i^1 X_j^1}{r^5} - \frac{X_i^1 X_j^1}{r^3} \delta_{ij} \right) \end{aligned}$$

$$\text{NR: } \sum_j \frac{x_j^{12}}{r^3} = \frac{r^{12}}{r^3} = \frac{r^{12} r^2}{r^5} = \frac{r^{12}}{r^5} \sum_i x_i x_i = \frac{1}{r^5} \sum_i \sum_j x_i x_j r^{12} \delta_{ij}$$

$$\Downarrow \Phi = \frac{1}{4\pi\epsilon_0} \int_{(\vec{x}') \in R} \frac{\rho(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x'$$

$$\Phi = \frac{1}{4\pi\epsilon_0} \left( \frac{Q}{r} + \frac{1}{r^3} \sum_i x_i p_i + \frac{1}{r^5} \sum_i \sum_j x_i x_j Q_{ij} + \dots \right)$$

Multipolentwicklung

$$\text{mit } Q = \int \rho(\vec{x}') d^3x' \quad \hat{=} \text{ Gesamtladung}$$

$$p_i = \int \rho(\vec{x}') x_i' d^3x' \quad \hat{=} \text{ Dipolmoment}$$

$$Q_{ij} = \int \rho(\vec{x}') (3 x_i' x_j' - \delta_{ij} r'^2) d^3x' \quad \hat{=} \text{ Quadrupolmoment}$$

$$\vdots$$

Eigenschaften von  $Q_{ij}$ :

- $Q_{ij} = Q_{ji}$
- $\sum_i Q_{ii} \equiv \text{Sp}(Q) = 0$ 

$$= \sum_i \int \mathcal{L} (3X_i'^2 - \delta_{ii} r'^2) d^3r$$

$$= \int \mathcal{L} (3r'^2 - 3r'^2) d^3r = 0$$

$\Rightarrow Q_{ij}$  hat 5 unabhängige Elemente (9 - 3 - 1)

Beispiele: i) Punktladung am Ursprung  $\mathcal{L}(\vec{x}') = q \mathcal{L}(\vec{x}'^2)$

$\Rightarrow \vec{E} = \frac{q}{r^2} = \text{Monopol}$ ,  $P_i = Q_{ij} = 0$

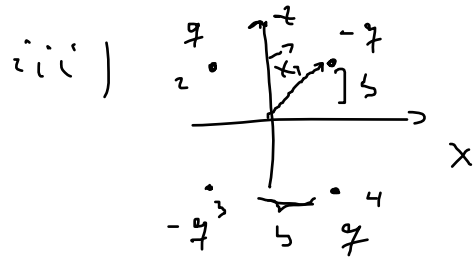
ii)  $\begin{matrix} z \\ \uparrow \\ b \{ \cdot \\ \uparrow +q \\ b \{ \cdot \\ \downarrow -q \end{matrix}$   $\rho(\vec{x}') = q \left[ \delta(\vec{x}' - \vec{L}) - \delta(\vec{x}' + \vec{L}) \right], \quad \vec{L} = b \vec{e}_z$

$$Q = \int \rho(\vec{x}') d^3x' = q - q = 0$$

$$P_i = \int \rho(\vec{x}') x'_i d^3x' = q(b_i - (-b_i)) = 2q b_i$$

$$\vec{P} = 2q \vec{L}, \quad q \vec{d}, \quad \vec{d} = 2\vec{L}$$

$$\begin{aligned} Q_{ij} &= \int \rho(\vec{x}') \cdot (3x'_i x'_j - \delta_{ij} r'^2) d^3x' \\ &= q \int (\delta(\vec{x}' - \vec{L}) - \delta(\vec{x}' + \vec{L})) (3x'_i x'_j - \delta_{ij} r'^2) d^3x' \\ &= q \left( 3 \cancel{b_i b_j} - \delta_{ij} \cancel{L^2} - 3(-b_i)(-b_j) + \delta_{ij} (-\cancel{L})^2 \right) = 0 \end{aligned}$$



$$S(\vec{x}) = -q \left[ \rho(\vec{x} - \vec{x}_1) - \rho(\vec{x} - \vec{x}_2) + \rho(\vec{x} - \vec{x}_3) - \rho(\vec{x} - \vec{x}_4) \right]$$

$$\vec{x}_1 = b(\vec{e}_x + \vec{e}_z)$$

$$\vec{x}_2 = b(-\vec{e}_x + \vec{e}_z)$$

$$\vec{x}_3 = b(-\vec{e}_x - \vec{e}_z)$$

$$\vec{x}_4 = b(\vec{e}_x - \vec{e}_z)$$

$$Q = 0, \quad P_i = -q \int \left[ \rho(\vec{x} - \vec{x}_i) - \rho(\dots) \right] x_i' d^3x'$$

$$P_y = 0, \quad P_x = -q b(1 - 1 - 1 + 1) = 0, \quad P_z = -q b(1 + 1 - 1 - 1) = 0$$

$$Q_{ij} = \int S(\vec{x}') \left( 3x_i' x_j' - \delta_{ij} \underbrace{r'^2}_{2b^2} \right) d^3x'$$

$$Q_{11} = -q b^2 (1 - 1 + 1 - 1) = 0, \quad Q_{22} = 0 = Q_{33}$$

$$Q_{12} = 0 = Q_{23}, \quad Q_{13} = -3q b^2 (1 - (-1)1 + (-1)(-1) - (1)(1)) = \underline{\underline{-12q b^4}}$$





lin  $c_m$  mit  $\Delta$  existieren,  $f_{m,0} = 0$   
 $r' \rightarrow 0$

$$\downarrow \quad \downarrow \\ |\vec{X} - \vec{X}'| = \sum_c \sum_m \underset{\substack{\uparrow \\ \text{konstant}}}{d_{cm}} \frac{r'^l}{r^{l+m}} Y_{cm}^*(\theta', \varphi') Y_{cm}(\theta, \varphi)$$

Satz: Additionstheoreme für Kugelfunktionen

$$P_l(\cos \delta) = \frac{4\pi}{2l+1} \sum_{m=-l}^l Y_{lm}^*(\theta', \varphi') Y_{lm}(\theta, \varphi)$$

$$\gamma = \angle(\vec{X}, \vec{X}') \quad , \text{d.h.} \quad \vec{X} \cdot \vec{X}' = |\vec{X}| |\vec{X}'| \cos \delta(\theta, \varphi, \theta', \varphi')$$

Beleg: [Jackson]

$$\cos \gamma = 1, \quad \vec{x} \parallel \vec{x}', \quad \theta = \theta', \quad \rho = \rho'$$

$$P_e(1) = \frac{4\pi}{2e+1} \sum_m Y_{em}^*(\theta, \varphi) Y_{em}(\theta, \varphi)$$

$$P_e(s) := \frac{1}{2^e e!} \left. \frac{d}{ds^e} (s^2 - 1)^e \right|_{s=1} = \frac{e!}{2^e e!} \left( \underbrace{\frac{d}{ds} (s^2 - 1)}_{2s} \right)^e \Big|_{s=1} = \frac{e! 2^e}{2^e e!} = 1$$

$$\frac{1}{|\vec{x} - \vec{x}'|} = \frac{1}{\sqrt{r^2 + r'^2 - 2rr' \cos \gamma}} \stackrel{\uparrow}{=} \frac{1}{\sqrt{(r-r')^2}} = \frac{1}{r-r'} = \frac{1}{r} \frac{1}{1-\frac{r'}{r}}$$

$\cos \gamma = 1$

$$= \frac{1}{r} \underbrace{\sum_{e=0}^{\infty} \left(\frac{r'}{r}\right)^e}_{\frac{1}{1-\frac{r'}{r}}} = \sum_{e=0}^{\infty} \frac{r'^e}{r^{e+1}} = \sum_e \sum_m \text{den} \frac{r'^e}{r^{e+1}} Y_{em}^*(\theta, \varphi) Y_{em}(\theta, \varphi)$$

geo. Reihe

$$\Rightarrow \sum_m \text{den} Y_{em}^*(\theta, \varphi) Y_{em}(\theta, \varphi) = 1$$

$$\Rightarrow \boxed{\text{den} = \frac{4\pi}{2e+1}}$$

$$\frac{1}{|\vec{x} - \vec{x}'|} \stackrel{|\vec{x}| \gg R}{=} \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \frac{4\pi}{2\ell+1} \frac{r'^{\ell}}{r^{\ell+1}} Y_{\ell m}^*(\theta', \varphi') Y_{\ell m}(\theta, \varphi)$$

$$\Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3x'$$

$$= \frac{1}{4\pi\epsilon_0} \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \sqrt{\frac{4\pi}{2\ell+1}} \frac{q_{\ell m}}{r^{\ell+1}} Y_{\ell m}(\theta, \varphi)$$

$$q_{\ell m} = \sqrt{\frac{4\pi}{2\ell+1}} \int \rho(\vec{x}') r'^{\ell} Y_{\ell m}^*(\theta', \varphi') d^3x'$$

$$q_{00} = \sqrt{4\pi} \int P(\vec{x}') \underbrace{Y_{00}^*}_{\frac{1}{4\pi}} d^3x' = \frac{\Phi}{\sqrt{4\pi}}$$

$$\Phi = \frac{1}{4\pi \epsilon_0} \sqrt{4\pi} \frac{q_{00}}{\sqrt{}} = \frac{1}{4\pi \epsilon_0} \frac{\Phi}{\sqrt{}}$$

$$l=1, m=0, \pm 1 \quad q_{10} = \sqrt{\frac{4\pi}{3}} \int P(\vec{x}') \underbrace{Y_{10}^*}_{\sqrt{\frac{3}{4\pi}} \cos\theta} d^3x' = \int P(\vec{x}') z' d^3x = P_z$$

and 
$$P_x = \frac{1}{\sqrt{2}} (q_{1,-1} - q_{1,1})$$

$$P_z = \frac{1}{\sqrt{2i}} (q_{1,1} + q_{1,-1})$$

$l=2, m=0, \pm 1, \pm 2 \Rightarrow 5$  Quadrupolmomente  $q_{2m}$   
hängen mit  $Q_{ij}$  zusammen