

# Lagrangian Formalismus der E-Dynamik

Mechanik:  $S = \int L(q, \dot{q}, t) dt$

$\nearrow$  Wirkung  
 $\nearrow$  Lagrange fkt.

$$\delta S = 0 \Rightarrow \frac{\partial L}{\partial q} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} = 0 \quad \text{E-L-Gl.}$$

Verallgemeinerung auf Felder

Mechanik:  $q(t) \rightarrow \underbrace{\Phi(\vec{x}, t)}_{x^m} \Rightarrow \mathcal{L}[\Phi, \dot{\Phi}]$

$$\bullet \quad t \rightarrow \vec{x}, \quad t \Rightarrow x^\mu$$

$$S[\Phi, \partial\Phi] = \iiint \mathcal{L}[\Phi, \partial\Phi] \underbrace{d^3x dt}_{\equiv d^4x}$$

$$\delta S = 0 = \int \left\{ \underbrace{\frac{\partial \mathcal{L}}{\partial \Phi}}_{\text{functional}} \delta \Phi + \sum_{\mu} \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \Phi)} \underbrace{\delta (\partial_{\mu} \Phi)}_{\partial_{\mu} \delta \Phi} \right\} d^4x$$

functional  
ableitung

$$= \int \left\{ \frac{\partial \mathcal{L}}{\partial \Phi} \delta \Phi - \sum_{\mu} \left( \partial_{\mu} \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \Phi)} \right) \delta \Phi \right\} d^4x + \underbrace{\sum_{\mu} \int \partial_{\mu} \left( \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \Phi)} \delta \Phi \right)}_{d^4x}$$

Annahme:  $\delta \Phi|_{\text{Rand}} = 0$

$$\frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \Phi)} \delta \Phi \Big|_{\text{Rand}} = 0$$

$$0 = \int \underbrace{\left\{ \frac{\partial \mathcal{L}}{\partial \Phi} - \sum_{\mu=0}^3 \partial_{\mu} \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \Phi)} \right\}}_{\Rightarrow = 0} \delta \Phi \, d^4 x$$

E-L Gleichung eines Feldtheorie:

$$\frac{\partial \mathcal{L}}{\partial \Phi} - \sum_{\mu=0}^3 \partial_{\mu} \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \Phi)} = 0$$

$$x^{\mu} = (x^0, x, y, z)$$

"   
 = t

analog zu

$$\frac{\partial \mathcal{L}}{\partial q} - \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}} = 0$$

Beispiel:  $\mathcal{L} = \frac{1}{2} \sum_K \sum_\nu (\partial_K \Phi) (\partial_\nu \Phi) \gamma^{K\nu} = \frac{1}{2} \sum_K (\partial_K \Phi) (\partial^\nu \Phi)$

$= \delta^K \Phi$

$$\frac{\partial \mathcal{L}}{\partial \Phi} = 0$$

$$\frac{\partial \mathcal{L}}{\partial (\partial_\mu \Phi)} = \frac{1}{2} \sum_K \sum_\nu \left( \underbrace{\frac{\partial (\partial_K \Phi)}{\partial (\partial_\mu \Phi)}}_{\delta_R^\mu} \partial_\nu \Phi \gamma^{K\nu} + \gamma^{K\nu} \partial_\nu \Phi \frac{\partial (\partial_\nu \Phi)}{\partial (\partial_\mu \Phi)} \right)$$

$$= \frac{1}{2} \sum_K \sum_\nu \left( \underbrace{\frac{\partial (\partial_K \Phi)}{\partial (\partial_\mu \Phi)}}_{\delta_R^\mu} \partial_\nu \Phi \gamma^{K\nu} + (\partial_\nu \Phi) \underbrace{\frac{\partial (\partial_\nu \Phi)}{\partial (\partial_\mu \Phi)}}_{\delta_\nu^\mu} \gamma^{K\nu} \right)$$

$$\begin{aligned} \frac{\partial \mathcal{L}}{\partial (\partial_\mu \Phi)} &= \frac{1}{2} \sum_{\nu} \partial_\nu \Phi \eta^{\mu\nu} + \frac{1}{2} \sum_{\nu} (\partial_\nu \Phi) \eta^{\mu\nu} \\ &= \sum_{\nu} (\partial_\nu \Phi) \eta^{\mu\nu} \end{aligned}$$

$$\begin{aligned} \sum_{\mu} \partial_\mu \frac{\partial \mathcal{L}}{\partial (\partial_\mu \Phi)} &= \sum_{\mu} \sum_{\nu} (\partial_\mu \partial_\nu \Phi) \eta^{\mu\nu} = \underbrace{\sum_{\mu} \sum_{\nu} \eta^{\mu\nu} \partial_\mu \partial_\nu \Phi}_{\square} \\ &= \square \Phi \end{aligned}$$

$\Rightarrow$  E-L Gl.

$$\square \Phi = 0$$

well known

$$\rightarrow \mathcal{L} = \frac{1}{2} \sum_{\mu} \sum_{\nu} (\partial_{\mu} \Phi)(\partial_{\nu} \Phi) \eta^{\mu\nu}$$

ist Lagrangian für Wellenfelder  $\partial^2 \Phi = 0$

(analog: der Feder-L:  $L = \frac{1}{2} m \dot{x}^2$ )

$$EL: \ddot{x} = 0$$

Wicht. Schl.: ED

$$\mathcal{L} = -\frac{1}{4} \sum_{\mu} \sum_{\nu} F_{\mu\nu} F^{\mu\nu} - \mu_0 \underbrace{\sum_{\mu} A^{\mu} j_{\mu}}_{\sum_{\mu\nu} A_{\mu} j_{\nu} \gamma^{\mu\nu}}$$

$$F_{\mu\nu} = \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu} \quad \text{Feldstärke tensor}$$

$$A_{\mu} = (\vec{A}, \vec{\Phi})$$

E-L Gl.:

$$\frac{\partial \mathcal{L}}{\partial A_{\mu}} - \sum_{\lambda=0}^3 \partial_{\lambda} \frac{\partial \mathcal{L}}{\partial (\partial_{\lambda} A_{\mu})} = 0 \quad \begin{array}{l} \text{4 Gl.} \\ k=0,1,2,3 \end{array}$$

$$\frac{\partial \mathcal{L}}{\partial A_{\mu}} = -\mu_0 \sum_{\nu} \sum_{\rho} \gamma^{\nu\rho} j_{\nu} \underbrace{\frac{\partial A_{\nu}}{\partial A_{\mu}}}_{\delta_{\nu}^{\mu}} = -\mu_0 \sum_{\nu} \gamma^{\mu\nu} j_{\nu} = -\mu_0 j^{\mu}$$

$$\frac{\partial \mathcal{L}}{\partial(\partial_\lambda A_\mu)} = -\frac{1}{4} \sum_\mu \sum_\nu \left\{ \frac{\partial F_{\mu\nu}}{\partial(\partial_\lambda A_\mu)} F^{\mu\nu} + F_{\mu\nu} \frac{\partial F^{\mu\nu}}{\partial(\partial_\lambda A_\mu)} \right\}$$

$$= -\frac{1}{2} \sum_\mu \sum_\nu \underbrace{\frac{\partial F_{\mu\nu}}{\partial(\partial_\lambda A_\mu)}}_{\text{}} F^{\mu\nu}$$

$$\hookrightarrow \frac{\partial F_{\mu\nu}}{\partial(\partial_\lambda A_\mu)} = \frac{\partial(\partial_\mu A_\nu)}{\partial(\partial_\lambda A_\mu)} - \frac{\partial(\partial_\nu A_\mu)}{\partial(\partial_\lambda A_\mu)} = \delta_\mu^\lambda \delta_\nu^{\mu} - \delta_\nu^\lambda \delta_\mu^{\mu}$$

$$\uparrow \frac{\partial \mathcal{L}}{\partial(\partial_\lambda A_\mu)} = -\frac{1}{2} \sum_\mu \sum_\nu (\delta_\mu^\lambda \delta_\nu^{\mu} - \delta_\nu^\lambda \delta_\mu^{\mu}) F^{\mu\nu}$$

$$= -\frac{1}{2} \begin{pmatrix} F^{12} & -F^{21} \\ \underbrace{-F^{12}} \end{pmatrix} = -F^{12}$$



$$\downarrow \frac{\partial \mathcal{L}}{\partial A_{\lambda k}} - \sum_{\lambda} \partial_{\lambda} \frac{\partial \mathcal{L}}{\partial (\partial_{\lambda} A_{\lambda k})} = 0$$

$$\Rightarrow -\mu_0 j^k + \sum_{\lambda} \partial_{\lambda} F^{\lambda k} = 0$$

$$\Rightarrow \boxed{\sum_{\lambda} \partial_{\lambda} F^{\lambda k} = \mu_0 j^k}$$

inhomogen  
Maxwell-Gleichung  
(gesehen in letztem Vorlesung)

$$\Rightarrow \boxed{\mathcal{L} = -\frac{1}{4} \sum_{\mu} \sum_{\nu} F_{\mu\nu} F^{\mu\nu} - \mu_0 \sum_{\mu} A^{\mu} j_{\mu}}$$

ist Lagrangendichte für E-Dynamik

homogene Maxwell-Gleichungen sind keine E-L-Gleichungen

Sondern Bianchi-Identitäten:

$$\partial_\lambda F_{\mu\nu} + \partial_\mu F_{\nu\lambda} + \partial_\nu F_{\lambda\mu} = 0 \quad (\text{letzte Vorlesung})$$

$$= \partial_\lambda (\cancel{\partial_\mu A_\nu} - \cancel{\partial_\nu A_\mu}) + \partial_\mu (\cancel{\partial_\nu A_\lambda} - \cancel{\partial_\lambda A_\nu}) + \partial_\nu (\cancel{\partial_\lambda A_\mu} - \cancel{\partial_\mu A_\lambda}) = 0$$

ist identisch erfüllt, also eine Identität

## Beweis

$$(i) \sum_K \partial_K \left( \sum_\lambda \partial_\lambda F^{\lambda K} \right) = \sum_K \sum_\lambda \partial_K \partial_\lambda F^{\lambda K} = 0$$

$$= \mu_0 \sum_{12} \partial_{12} j^K$$

$$\Rightarrow \boxed{\sum_K \partial_{12} j^K = 0}$$

Kontinuitätsgleichung

$$\vec{\nabla} \cdot \vec{j} + \frac{\partial \rho}{\partial t} = 0$$

ii)  $\mathcal{L}$  ist Lorentz invariant

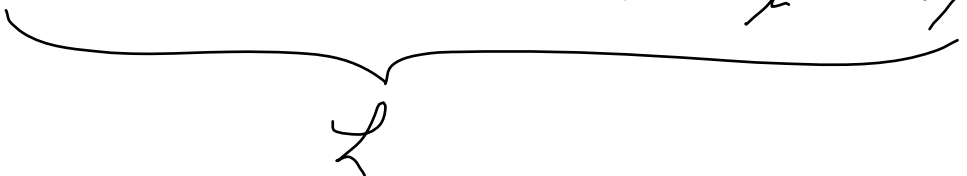
iii)  $\mathcal{L}$  ist eil invariant

Eichtransf.  $A_\mu \rightarrow A'_\mu = A_\mu - \partial_\mu \lambda$

lok. Vorl.  $F_{\mu\nu} \rightarrow F'_{\mu\nu} = F_{\mu\nu}$

$$\mathcal{L} \rightarrow \mathcal{L}' = -\frac{1}{4} \sum_{\mu} \sum_{\nu} F'_{\mu\nu} F'^{\mu\nu} - \mu_0 \sum_{\mu} A'^{\mu} j_{\mu}$$

$$= -\frac{1}{4} \sum_{\mu} \sum_{\nu} F_{\mu\nu} F^{\mu\nu} - \mu_0 \sum_{\mu} A^{\mu} j_{\mu} + \mu \sum_{\mu} (\partial^{\mu} \lambda) j_{\mu}$$



$$= \mathcal{L} + \mu_0 \sum_{\mu} (\partial^{\mu} \lambda) j_{\mu}$$

d.h.  $\mathcal{L}$  ist nicht eichinvariant

$$\begin{aligned}
\text{aber } S &\rightarrow S' = \int \mathcal{L}' d^4x \\
&= \underbrace{\int \mathcal{L} d^4x}_S + \mu_0 \sum_{\mu} \int (\partial^{\mu} \lambda) \hat{j}_{\mu} d^4x \\
&= S - \mu_0 \sum_{\mu} \int \lambda \underbrace{(\partial^{\mu} \hat{j}_{\mu})}_{\substack{=0 \text{ Kontinuitätsgl.}}} d^4x \\
&\quad + \mu_0 \sum_{\mu} \underbrace{\lambda \hat{j}_{\mu}}_{=0} \Big|_{\partial=1} \\
&= S
\end{aligned}$$

$\Rightarrow S$  ist eil invariant  $\Rightarrow$  E-L. Ge<sup>Sind</sup> weil invariant !

Klausur

Termin: 19. 2., 10-12 Uhr, HS II

Hilfsmittel: - 2 Din A4 Blätter Land beschriften  
 - kein Handy, kein Taschenrechner, kein Computer  
 Papier mitbringen

Vorklausur: • a) how many Trefts / Konten  
 • Depreciation 17. 2. 10<sup>30</sup> HS III

Vorbereitung: Blatt 1-3 komplett, (4/1,2), (5/1,2,4)  
 (6/1,3), (7/1), (8/1,3), (9/1,2), (10/1,3,4)  
 (11/2,3), (12/1,2,4), (13/2,3)

(ohne Beweise)